

# General Relativity without paradigm of space-time covariance: sensible quantum gravity and resolution of the problem of time

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The framework of a theory of gravity from quantum to classical regimes is presented. Paradigm shift from full space-time covariance to only spatial diffeomorphism invariance, together with clean decomposition of the canonical structure, yield transparent physical dynamics and a resolution of the problem of time. The deep divide between quantum mechanics and conventional canonical formulations of quantum gravity is overcome with a Schrodinger equation for quantum geometrodynamics which describes evolution in intrinsic time. Unitary time development with gauge-invariant temporal ordering is also viable. All Kuchar observables become physical; and classical space-time, with direct correlation between its proper times and intrinsic time intervals, emerges from constructive interference. The framework not only yields a physical Hamiltonian for Einstein's theory, but also prompts natural extensions and improvements towards a well-behaved quantum theory of gravity.

PACS numbers: 04.60.-m, 04.60.Ds

Covariance of space and time has been crucial to Einstein's theory of General Relativity (GR); but the assumption of full 4-dimensional (4d) diffeomorphism symmetry also entails a number of technical and conceptual difficulties. On the technical front, Einstein's theory fails to be renormalizable as a perturbative quantum field theory; and GR remains incomplete as a quantum theory, despite many recent advances. At the conceptual level, the Hamiltonian constraint which is believed to dictate the dynamics has a dual role as a generator of symmetry, with arbitrary lapse function as Lagrange multiplier. Full 4d invariance with local Hamiltonian constraint and its consequent baggage of arbitrary lapse and gauged histories is hard to reconcile with the physical reality of time.

Canonical quantum gravity has the guise of a formalism "frozen in time". Quantum states do not evolve in coordinate "time", and  $dx^0$  is merely one component of space-time displacement, with no invariant physical meaning in a theory with full space-time covariance. Yet a notion of time is needed to make sense of quantum mechanical interpretations: both for the discussion of dynamics and evolution, and for the interpretation of probabilities which are normalized at particular instants of time. Even if a suitable degree of freedom (d.o.f.) is chosen as the intrinsic time (as many have advocated), a successful quantum theory will still need to reconcile a Klein-Gordon type Wheeler-DeWitt (WDW) equation[1, 2] quadratic in momenta with the positivity of "probabilities". A resolution of the "problem of time"[3] cannot be deemed complete if it fails to account for the intuitive physical reality of time and does not provide a satisfactory correlation between intrinsic time development in quantum dynamics and the passage of time in classical space-times.

At the deeper level of symmetries, the constraints of GR satisfy the Dirac algebra[4]. However, closer inspection reveals it corresponds to 4d diffeomorphisms only modulo equations of motions (EOM). Off-shell, full 4d Lie derivatives involve time derivatives; these cannot be generated by constraints which depend only on the spatial metric  $q_{ij}$  and its conjugate momentum  $\tilde{\pi}^{ij}$ . In Einstein's theory, the EOM, among other things, relate  $\dot{q}_{ij}$  to  $\tilde{\pi}^{ij}$  and the constraints then generate 4d diffeomorphisms on-shell. This is a clue that, without the help of EOM, a *quantum* theory of GR cannot, and need not, enforce full 4d space-time covariance *off-shell*.

A key obstacle to the viability of GR as a perturbative quantum field theory lies in the conflict between unitarity and space-time general covariance: renormalizability can be attained with higher-order curvature terms, but space-time covariance requires time as well as spatial derivatives of the same (higher) order, thus compromising unitarity. Horava relinquished full 4d symmetry and achieved power-counting renormalizable modifications of GR with higher-order spatial curvature terms[5]. In loop quantum gravity, the non-perturbative master constraint program[6] seeks representations not of the Dirac algebra, but of the master constraint algebra which has the advantages of having structure *constants* (rather than functions), and of decoupling the equivalent quantum Hamiltonian constraint from spatial diffeomorphism generators  $H_i$ . Ref.[7] consistently realized Horava gravity theories as canonical theories with first class master constraint algebra. This not only removes the canonical inconsistencies of projectable Horava theory, but also captures the essence of the theory in retaining spatial diffeomorphisms as the *only* local gauge symmetries. With full covariance, observables,  $O$ , of Einstein's theory are required to commute both with  $H_i$  and  $H$ , leading to the demand (above and beyond usual canonical rules) of  $\{O, \{O, M\}\}_{M=0} = 0$ [6]. In contradistinction, in this work a theory of gravity appositely formulated with master constraint marks a real paradigm shift in the symmetry, from full 4d general coordinate invariance to invariance only with respect to spatial diffeomorphisms. Enforcing  $H(x) = 0$  through the master constraint effectively removes it from one of its dual roles since  $M$  generates no extra symmetry, and determines only dynamics.

A theory of geometrodynamics, with  $(q_{ij}, \tilde{\pi}^{ij})$  as fundamental variables, is bequeathed with a number of remarkable features: positivity of the spatial metric guarantees space-like separation between any two points on the initial value hypersurface (and allows, to quote Ref.[2], "a notion of 'simultaneity' and a common moment of a rudimentary 'time'"). In Arnowitt-Deser-Misner (ADM)[8] description of the space-time metric, these are labeled as constant- $x^0$  hypersurfaces. Quantum states however do not depend on  $x^0$ ; the link between intrinsic and coordinate times which will be revealed is more subtle. Decomposition of  $q_{ij} = q^{\frac{1}{3}} \bar{q}_{ij}$  into determinant and unimodular factors results in clean separation of canonical pairs. The generic ultra-local DeWitt supermetric[1] with deformation parameter  $\lambda$  has signature  $(\text{sgn}[\frac{1}{3} - \lambda], +, +, +, +, +)$ , and the single negative eigenvalue for  $\lambda > \frac{1}{3}$  corresponds to the  $\delta \ln q$  mode. It is the apposite choice (as subsequent discussions will show) for intrinsic time interval in quantum gravity. Although  $q$  is a scalar density (a criticism which has been used to disqualify it from the role of time variable),  $\delta \ln q = \frac{\delta q}{q}$  is a scalar. Even in Galilean-Newtonian physics, it is *time interval*, and not absolute time, which is physical.

With the preceding concerns and observations in mind, the framework for a theory of gravity, from quantum to classical regimes, will be presented wherein several factors collude to result in a minimal and compelling formulation. The WDW constraint permits factorization into  $(\beta\pi - \bar{H})(\beta\pi + \bar{H}) = 0$ ; and the vanishing of only one of the factors is sufficient for capturing the classical content of GR (with  $\beta^2 = \frac{1}{6}$ ). In the master constraint formulation wherein the Hamiltonian constraint of GR is freed from its role of generator of "time"-development with arbitrary lapse, the commencing action is  $\int [\tilde{\pi}^{ij} \dot{q}_{ij} - N^i H_i] dt - \int m(t) M dt$ , with  $M := \int (\beta\pi + \bar{H})^2 / \sqrt{q} = 0$ . The resultant

first class constraint algebra,  $\{M, M\} = 0, \{H_i[N^i], M\} = 0, \{H_i[N^i], H_j[N'^j]\} = H_i[\mathcal{L}_{\vec{N}'} N'^i]$ , exhibits only spatial diffeomorphism gauge symmetry, both on- and off-shell. That  $\pi$ , the trace of the momentum, is conjugate to  $\ln q^{\frac{1}{3}}$  leads to a quantum theory described by a Schrodinger equation *first order in intrinsic time*, with consequent positive semi-definite probability density  $\Psi^\dagger \Psi$ . Formulation in the Heisenberg picture and gauge-invariant temporal-ordering of the evolution operator are also viable. There is the additional promise of a well-behaved unitary quantum theory of gravity if the Hamiltonian generating intrinsic time development,  $\mathcal{H}_{phys.}$ , can be made real, bounded from below, and also renormalizable by including suitable higher-order spatial curvature terms (with GR recovered in the limit of low curvatures). The resultant semi-classical Hamilton-Jacobi (HJ) equation is also first order, with the implication of completeness[9] (its integral solution provides a complete set of gauge-invariant integration constants of motion). Classical space-time emerges from constructive interference of quantum wave functions[10]. The physical content of GR is regained from a theory with  $\bar{H}/\beta$  generating intrinsic  $\delta \ln q^{\frac{1}{3}}$  time translations, subject to only spatial diffeomorphism invariance. Furthermore, the emergent classical space-times have ADM metrics with proper times (for vanishing shifts) directly correlated to intrinsic time intervals by the explicit relation  $d\tau^2 = [\frac{\delta \ln q}{(12\beta\kappa H/\sqrt{q})}]^2$ .

## THEORY OF GRAVITY WITHOUT PARADIGM OF FULL SPACE-TIME COVARIANCE

In the initial value problem, York[12] explored a conformal decomposition of the spatial metric  $q_{ij} = \phi \bar{q}_{ij}$ , of which  $\phi = q^{\frac{1}{3}}$  with unimodular  $\bar{q}_{ij}$  is a special case ( $q := \det[q_{ij}]$ ). However, in focusing on the generic conformal factor  $\phi$ , the “miracle” of  $\phi = q^{\frac{1}{3}}$  was not fully revealed. Among other advantages, it results in clean separation of  $(\ln q^{\frac{1}{3}}, \pi)$  from other canonical pairs, with the symplectic potential,

$$\int \tilde{\pi}^{ij} \delta q_{ij} = \int \tilde{\pi}^{ij} \delta \bar{q}_{ij} + \pi \delta \ln q^{\frac{1}{3}}, \quad (1)$$

wherein  $\pi := q_{ij} \tilde{\pi}^{ij}$  and  $\tilde{\pi}^{ij} := q^{\frac{1}{3}}[\tilde{\pi}^{ij} - \frac{q^{ij}}{3}\pi]$ . Thus the only non-trivial Poisson brackets are  $\{\bar{q}_{kl}(x), \tilde{\pi}^{ij}(x')\} = P_{kl}^{ij} \delta(x, x')$ , and  $\{\ln q^{\frac{1}{3}}(x), \pi(x')\} = \delta(x, x')$ ; with the trace-free projector,  $P_{kl}^{ij} := \frac{1}{2}(\delta_k^i \delta_l^j + \delta_l^i \delta_k^j) - \frac{1}{3}q^{ij} \bar{q}_{kl}$ . This separation carries over to the quantum theory, and permits a d.o.f., separate from the others, to be identified as the carrier of temporal information. An intrinsic clock is in fact not tied to 4d general covariance. A generic ultra-local DeWitt supermetric[1] compatible with spatial diffeomorphism invariance,  $G_{ijkl} = \frac{1}{2}(q_{ik}q_{jl} + q_{il}q_{jk}) - \frac{\lambda}{3\lambda-1}q_{ij}q_{kl}$ , comes equipped with intrinsic temporal intervals  $\delta \ln q^{\frac{1}{3}}$  provided  $\lambda > \frac{1}{3}$ . Crucially it is also  $\ln q^{\frac{1}{3}}$  which can be so neatly separated from the rest. With  $\beta^2 := \frac{1}{3(3\lambda-1)}$ , the Hamiltonian constraint of a theory of geometrodynamics quadratic in momenta and with ultralocal supermetric is generically,

$$0 = \frac{\sqrt{q}}{2\kappa} H = G_{ijkl} \tilde{\pi}^{ij} \tilde{\pi}^{kl} + V(q_{ij}) = -(\beta\pi - \bar{H})(\beta\pi + \bar{H});$$

$$\bar{H}(\tilde{\pi}^{ij}, \bar{q}_{ij}, q) = \sqrt{\bar{G}_{ijkl} \tilde{\pi}^{ij} \tilde{\pi}^{kl} + V(\bar{q}_{ij}, q)} = \sqrt{\frac{1}{2}[\bar{q}_{ik}\bar{q}_{jl} + \bar{q}_{il}\bar{q}_{jk}] \tilde{\pi}^{ij} \tilde{\pi}^{kl} + V(q_{ij})}. \quad (2)$$

Einstein’s GR ( $\lambda = 1$  and  $V(\bar{q}_{ij}, q) = -\frac{q}{(2\kappa)^2}[R - 2\Lambda_{eff}]$ ) is a particular realization of a wider class of theories.

Several features are noteworthy: (i) As written, equation (2) is a local constraint, in addition to those of spatial diffeomorphism invariance  $H_i = 0$ . This leads to a quandary: if the constraint algebra is first class, then  $H$  consistently generates multi-fingered time translation *symmetry* as well as “physically” evolving the theory with respect to “time” (as manifested by its dual roles in Einstein’s theory with full general coordinate invariance); (ii) a master constraint formulation with  $M = \int H^2/\sqrt{q} = 0$  equivalently enforces the local content of (2). This permits  $H$  to determine dynamical evolution rather than generate symmetry.  $M$  also decouples from  $H_i$  in the first class constraint algebra, paving the road for quantization[6]. The result is a theory with only spatial diffeomorphism invariance; with physical dynamics dictated by  $H$ , but encoded in  $M$ . (iii) The total constraints of the theory generate only spatial diffeomorphisms since  $\{f(q_{ij}, \tilde{\pi}^{ij}), m(t)M + H_k[N^k]\}_{|M=0 \Leftrightarrow H=0} \approx \{f, H_k[N^k]\} = \mathcal{L}_{\vec{N}} f$ . Thus  $m(t)$  plays no role in the dynamics. Instead, true physical evolution can only be with respect to an intrinsic time extracted from the WDW constraint. As detailed above,  $\ln q^{\frac{1}{3}}$  is the preeminent choice. (iv) The factorization allows for the possibility that  $(\beta\pi + \bar{H}) = 0$  is sufficient to recover the dynamical content of GR. This is a breakthrough: the semiclassical HJ equation is first order in intrinsic time (since  $\pi$  is conjugate to  $\ln q^{\frac{1}{3}}$ ), with its consequence of completeness[9]; and quantum gravity will now be dictated by a corresponding Schrodinger equation first order in intrinsic time (with consequent positive semi-definite probability density at any instant of intrinsic time). This resolves the deep divide

between a quantum mechanical interpretation (in which *both* the notion of time and positive semi-definite probabilities are needed) and the usual Klein-Gordon type WDW equations second-order in intrinsic time (hence indefinite in “probabilities”). In general  $\beta = \pm\sqrt{\beta^2}$ , but the positive value for  $\beta$  is singled out to render the physical Hamiltonian  $\mathcal{H}_{phys.} = \int \bar{H}/\beta$  (which shall be discussed later) positive semidefinite. This implies  $\pi \leq 0$ ; which corresponds, in Friedmann-Robertson-Walker cosmological models, to an expanding universe.

A spatial diffeomorphism-invariant quantum theory with  $M := \int (\beta\pi + \bar{H})^2/\sqrt{q} = 0$  will translates into

$$[\beta\hat{\pi} + \bar{H}(\hat{\pi}^{ij}, \hat{q}_{ij}, \hat{q})]|\Psi\rangle = 0, \quad \hat{H}_i|\Psi\rangle = 0. \quad (3)$$

In the metric representation, canonical momenta are realized by  $\hat{\pi} = \frac{3\hbar}{i} \frac{\delta}{\delta \ln q}$ ,  $\hat{\pi}^{ij} = \frac{\hbar}{i} P_{lk}^{ij} \frac{\delta}{\delta \bar{q}_{lk}}$  which operate on  $\Psi[\bar{q}_{ij}, q]$ ; and the Schrodinger equation and HJ equation for semi-classical states  $Ce^{\frac{iS}{\hbar}}$  are respectively,

$$i\hbar \frac{\delta}{\delta \ln q} \Psi = \frac{\bar{H}(\hat{\pi}^{ij}, q_{ij})}{3\beta} \Psi, \quad \frac{\delta S}{\delta \ln q} = -\frac{\bar{H}(\bar{\pi}^{ij} = P_{kl}^{ij} \frac{\delta S}{\delta \bar{q}_{kl}}; q_{ij})}{3\beta}; \quad (4)$$

and  $\nabla_j \frac{\delta \Psi}{\delta \bar{q}_{ij}} = 0$  enforces spatial diffeomorphism symmetry. Behold the appearance of a true Hamiltonian  $\bar{H}/\beta$  generating evolution w.r.t. intrinsic time  $\delta \ln q^{\frac{1}{3}}(x)$ .

The master constraint formulation serves as a consistent canonical method to regain the physical content of Einstein’s theory, without paradigm of 4-covariance, from the usual starting point of canonical General Relativity with Hamiltonian  $\int (NH + N^i H_i)$ . (The analogy relativistic point particle mechanics is recounted in the Appendix.) The paradigm shift to spatial diffeomorphism invariance reveals the primacy of dynamics dictated by  $\bar{H}$  and the role of intrinsic time, and allows consistent extensions and improvements to Einstein’s theory.

### Emergence of classical spacetime

Many years ago Gerlach[10] demonstrated that classical space-time and its EOM can be recovered from the quantum theory through HJ theory and constructive interference. The first order HJ equation, which bridges quantum and classical regimes, has complete solution  $S = S^{(3)}\mathcal{G}; \alpha$  which depends on 3-geometry  $^{(3)}\mathcal{G}$  and integration constants (denoted generically here by  $\alpha$ ). Constructive interference with  $\mathcal{S}^{(3)}\mathcal{G}; \alpha + \delta\alpha = \mathcal{S}^{(3)}\mathcal{G}; \alpha$ ;  $\mathcal{S}^{(3)}\mathcal{G} + \delta^{(3)}\mathcal{G}; \alpha + \delta\alpha = \mathcal{S}^{(3)}\mathcal{G} + \delta^{(3)}\mathcal{G}; \alpha$  leads to  $\frac{\delta}{\delta \alpha} [\int \frac{\delta \mathcal{S}^{(3)}\mathcal{G}; \alpha}{\delta q_{ij}} \delta q_{ij}] = 0$ ; subject to constraints  $M = H_i = 0$ . With the momenta identified with  $\bar{\pi}^{ij}(\alpha) := \frac{\delta \mathcal{S}^{(3)}\mathcal{G}; \alpha}{\delta q_{ij}}$ , and Lagrange multipliers  $\delta m$  and  $\delta N^i$ , the requirement of constructive interference is equivalent to

$$\begin{aligned} 0 &= \frac{\delta}{\delta \alpha} [\int (\bar{\pi}^{ij} \delta q_{ij} + \delta N^i H_i) + \delta m M] \\ &= \frac{\delta}{\delta \alpha} [\int \pi \delta \ln q^{\frac{1}{3}} + \bar{\pi}^{ij} \delta \bar{q}_{ij} + \frac{q^{ij}}{3} \delta N_i \nabla_j \pi + q^{-\frac{1}{3}} \delta N_i \nabla_j \bar{\pi}^{ij}]. \end{aligned} \quad (5)$$

Happily, for master constraint theories, there is no  $\delta m$  contribution (since  $M = 0 \Leftrightarrow H = 0$  and  $M$  is quadratic in  $H$ ). Imposing  $\pi = -\bar{H}/\beta$ , integrating by parts, and bearing in mind  $\bar{H}(\bar{\pi}^{ij}(\alpha), q_{ij})$ , the resultant EOM is,

$$\frac{\delta \bar{q}_{ij}(x) - \mathcal{L}_{\bar{N}dt} \bar{q}_{ij}(x)}{\delta \ln q^{\frac{1}{3}}(y) - \mathcal{L}_{\bar{N}dt} \ln q^{\frac{1}{3}}(y)} = P_{ij}^{kl} \frac{\delta [\bar{H}(y)]}{\beta \delta \bar{\pi}^{kl}(x)} = \frac{\bar{G}_{ijmn} \bar{\pi}^{mn}}{\beta \bar{H}} \delta(x, y), \quad (6)$$

wherein  $\mathcal{L}_{\bar{N}}$  denotes Lie derivative and  $\delta \bar{N} =: \bar{N}dt$ . Proceeding as in Ref.[10], the other half of Hamilton’s equations,

$$\frac{\delta \bar{\pi}^{ij}(x) - \mathcal{L}_{\bar{N}dt} \bar{\pi}^{ij}(x)}{\delta \ln q^{\frac{1}{3}}(y) - \mathcal{L}_{\bar{N}dt} \ln q^{\frac{1}{3}}(y)} = -\frac{\delta [\bar{H}(y)/\beta]}{\delta \bar{q}_{ij}(x)}, \quad (7)$$

can be recovered. As predicted by (4),  $\bar{H}/\beta$  is the Hamiltonian for evolution of  $(\bar{q}_{ij}, \bar{\pi}^{ij})$  w.r.t.  $\ln q^{\frac{1}{3}}$ .

Although the derivation above bear similarities to Gerlach’s work, fundamental differences must be noted. In Ref.[10],  $\delta N =: Ndt$  (associated with local constraint  $H = 0$ ) will always contribute to the final EOM resulting in multi-fingered time with arbitrary lapse function. In contradistinction,  $\delta m$  contribution does not arise for a master constraint theory. This is part and parcel of the paradigm shift. Not only is unphysical time development with

arbitrary lapse function now evaded, the “defect” that  $M$  does not generate dynamical evolution w.r.t. coordinate time is redeemed at a much deeper level through physical evolution w.r.t. intrinsic time. Through (6) and  $\pi = -\bar{H}/\beta$ , the emergent ADM classical space-time has momentum

$$G_{ijkl}\tilde{\pi}^{kl} = \frac{\sqrt{q}}{4N\kappa} \left( \frac{dq_{ij}}{dt} - \mathcal{L}_{\bar{N}}q_{ij} \right), \quad Ndt := \frac{\delta \ln q^{\frac{1}{3}} - \mathcal{L}_{\bar{N}dt} \ln q^{\frac{1}{3}}}{(4\beta\kappa\bar{H}/\sqrt{q})}. \quad (8)$$

In conventional canonical formulation of Einstein’s GR, the EOM with *arbitrary* lapse is,

$$\frac{dq_{ij}}{dt} = \{q_{ij}, \int NH + N_i H^i\} = \frac{4N\kappa}{\sqrt{q}} G_{ijkl}\tilde{\pi}^{kl} + \mathcal{L}_{\bar{N}}q_{ij}. \quad (9)$$

The extrinsic curvature is related to  $\tilde{\pi}^{ij}$  by  $K_{ij} := \frac{1}{2N} \left( \frac{dq_{ij}}{dt} - \mathcal{L}_{\bar{N}}q_{ij} \right) = \frac{2\kappa}{\sqrt{q}} G_{ijkl}\tilde{\pi}^{kl}$ . Taking the trace yields

$$\frac{1}{3}Tr(K) = \frac{1}{2N} \left( \frac{\partial \ln q^{\frac{1}{3}}}{\partial t} - \mathcal{L}_{\bar{N}} \ln q^{\frac{1}{3}} \right) = \frac{2\kappa\beta}{\sqrt{q}} \bar{H}, \quad (10)$$

wherein the constraint  $(\beta\pi + \bar{H}) = 0$  has been used to arrive at the last step. (10) demonstrates that the lapse function and intrinsic time are precisely related (*a posteriori* by the EOM) by the same formula as in (8). For a theory with full 4d diffeomorphism invariance (such as Einstein’s GR with  $\beta^2 = 1/6$  and consistent Dirac algebra of constraints), this relation is an identity which does not compromise the arbitrariness of  $N$ . However, it reveals (even in Einstein’s GR) the physical meaning of the lapse function and its relation to the intrinsic time.

### Paradigm shift and resolution of the problem of time

Starting with only spatial diffeomorphism invariance and through constructive interference, equation (6) with physical evolution in intrinsic time generated by  $\bar{H}$  is obtained. This relates the momentum to coordinate time derivative of the metric precisely as in (8). It is thus possible to interpret the emergent classical space-time (which can generically be described with ADM metric) to possess extrinsic curvature which corresponds precisely to the derived lapse function displayed in (8). However, *only* the freedom of spatial diffeomorphism invariance is realized, as the *emergent* lapse is now completely described by the intrinsic time  $\ln q^{\frac{1}{3}}$  and  $\bar{N}$ . All EOM w.r.t coordinate time  $t$  generated by  $\int NH + N^i H_i$  in Einstein’s GR can be recovered from evolution w.r.t.  $\ln q^{\frac{1}{3}}$  and generated by  $\bar{H}/\beta$  iff  $N$  assumes the form of (8). All the previous observations lead to the central revelation: full 4d space-time covariance (with its consequent baggage of arbitrary lapse and gauged histories) is a red herring which obfuscates the physical reality of time, and all that is necessary to consistently capture the classical physical content of GR is a theory invariant only w.r.t. spatial diffeomorphisms accompanied by a master constraint which enforces the dynamical content. The paradigm shift points to a *complete resolution of the problem of time*, from quantum to classical GR: classical spacetime, with consistent lapse function and ADM metric,

$$ds^2 = - \left[ \frac{(\partial_t \ln q^{\frac{1}{3}} - \mathcal{L}_{\bar{N}} \ln q^{\frac{1}{3}}) dt}{4\beta\kappa(\bar{H}/\sqrt{q})} \right]^2 + q_{ij} [dx^i + N^i dt] [dx^j + N^j dt], \quad (11)$$

emerges from constructive interference of a spatial diffeomorphism invariant quantum theory with Schrodinger and HJ equations first order in intrinsic time development. Gratifying too are the correlations of classical proper time  $d\tau$  and quantum intrinsic time  $\ln q^{\frac{1}{3}}$  (through  $d\tau^2 = [\frac{\delta \ln q^{\frac{1}{3}}}{(4\beta\kappa\bar{H}/\sqrt{q})}]^2$  (for vanishing shifts)), and of Wheeler’s notion of ADM “simultaneity” to quantum simultaneity ( $dt = 0 \Rightarrow \delta \ln q = 0$  by equation (8)). In particular, by (9) and (10), proper time intervals measured by physical clocks in space-times which are solutions of Einstein’s equations always agree with the result of equation (11).

### Improvements to the quantum theory

The paradigm shift and framework presented prompt extensions and improvements to Einstein’s theory. Requirement of a real physical Hamiltonian  $\bar{H}$  compatible with spatial diffeomorphism symmetry suggests supplementing the

kinetic term in the square root with a positive semi-definite quadratic form, i.e.

$$\begin{aligned}\bar{H} &= \sqrt{\bar{G}_{ijkl}\bar{\pi}^{ij}\bar{\pi}^{kl} + \left[\frac{1}{2}(q_{ik}q_{jl} + q_{jk}q_{il}) + \gamma q_{ij}q_{kl}\right] \frac{\delta W}{\delta q_{ij}} \frac{\delta W}{\delta q_{kl}}} \\ &= \sqrt{[\bar{q}_{ik}\bar{q}_{jl} + \gamma \bar{q}_{ij}\bar{q}_{kl}](\bar{\pi}^{ij}\bar{\pi}^{kl} + q^{\frac{2}{3}} \frac{\delta W}{\delta q_{ij}} \frac{\delta W}{\delta q_{kl}})}.\end{aligned}\quad (12)$$

$\bar{H}$  is then real if  $\gamma > -\frac{1}{3}$ . To lowest order for perturbative power-counting renormalizability[5],  $W = \int [\sqrt{q}(aR - \Lambda) + g\tilde{\epsilon}^{ikj}(\Gamma_{im}^l \partial_j \Gamma_{kl}^m + \frac{2}{3}\Gamma_{im}^l \Gamma_{jn}^m \Gamma_{kl}^n)]$  (i.e. of the form of a 3-dimensional Einstein-Hilbert action with cosmological constant supplemented by a Chern-Simons action with dimensionless coupling constant  $g$ ). The potential of the form of Einstein's theory with cosmological constant is recovered at low curvatures. The effective value of  $\kappa$  and cosmological constant can thus be determined as  $\kappa = \frac{8\pi G}{c^3} = \sqrt{\frac{1}{2\pi^2 a \Lambda(1+3\gamma)}}$  and  $\Lambda_{eff} = \frac{3}{2}\kappa^2 \Lambda^2(1+3\gamma) = \frac{3\Lambda}{4a\pi^2}$  respectively. The possibility of having a new parameter  $\gamma$  in the potential (different from  $\lambda$  in the supermetric) was overlooked in previous works. Furthermore, positivity of  $\bar{H}^2$  (with  $\gamma > -\frac{1}{3}$ ) is correlated with *real*  $\kappa$  and *positive*  $\Lambda_{eff}$ . There is also the intriguing feature that the lowest classical energy of the physical Hamiltonian  $\bar{H}$  occurs when zero modes are present i.e.  $\gamma \rightarrow -\frac{1}{3}$ , leading, in this limit and for fixed  $\kappa$ , to  $\Lambda_{eff} \rightarrow 0$ . This, however, requires a thorough investigation of the renormalization group flow of  $\gamma$  and other parameters to deduce the exact behaviour of  $\Lambda_{eff}$  with physical energy scale, especially when matter and other forces are also taken into account. A slight generalization of (12) is to replace  $\frac{\delta W}{\delta q_{ij}}$  in the positive semidefinite quadratic form with  $\sqrt{q}(\Lambda' q^{ij} + a' R q^{ij} + b R^{ij} + g' C^{ij})$ , which is the most general symmetric second rank tensor (density) containing up to third derivatives of the spatial metric. The coupling constant  $g'$  associated with the Cotton-York tensor  $C^{ij}$  is dimensionless ( $\sqrt{q}C^{ij}$  is proportional to the functional derivative w.r.t. the spatial metric  $q_{ij}$  of the Chern-Simons term).

In Ref.[5], the Hamiltonian density with ‘‘detailed balance’’ is proportional to  $\frac{G_{ijkl}}{\sqrt{q}}(\pi^{ij}\pi^{kl} + \frac{\delta W}{\delta q_{ij}} \frac{\delta W}{\delta q_{kl}})$ . Without factoring out  $\pi$ , the negative mode in the full supermetric,  $G_{ijkl}$ , compromises the positivity of the kinetic term,  $\frac{G_{ijkl}}{\sqrt{q}}\pi^{ij}\pi^{kl}$ , in the theory. Formulations which use an extra scalar field[3], or variables other than  $\ln q^{\frac{1}{3}}$  as time to attain deparametrization will also be afflicted with the same problem. In contradistinction, with  $\ln q^{\frac{1}{3}}$  as intrinsic time,  $\pi$  is singled out and isolated as the conjugate variable, with the upshot that in the Schrodinger equation,  $\bar{H}$  (which does not contain  $\pi$ ) always has positive semidefinite kinetic term.

### Gauge-invariant global time, superspace dynamics, and temporal order

In our intrinsic time formulation, the Wheeler-DeWitt equation, which is a constraint which must be satisfied at each spatial point  $x$ , is replaced by a Schrodinger equation with  $\frac{\bar{H}(x)}{\beta}$  generating translations in  $\ln q^{\frac{1}{3}}(x)$  which is a Tomonaga-Schwinger[11] many-fingered time variable. The transcription from, apparently, many-fingered dynamics to evolution with respect to a gauge-invariant global time and the Heisenberg picture is, remarkably, unequivocal for 3-hypersurfaces which are compact Riemannian manifolds without boundary. Hodge decomposition of the 0-form  $\delta \ln q^{\frac{1}{3}}$  uniquely yields  $\delta \ln q^{\frac{1}{3}} = \delta h + \nabla_i \delta V^i$ , wherein  $\delta h$  is harmonic, *independent* of  $x$  and gauge-invariant, whereas  $\delta V^i$  can be gauged away because, fortuitously,  $\mathcal{L}_{\delta N^i} \ln q^{\frac{1}{3}} = \frac{2}{3} \nabla_i \delta N^i$ . This leads to the transcription,

$$i\hbar \frac{\delta \Psi}{\delta h} = \int i\hbar \left[ \frac{\delta \Psi}{\delta \ln q^{\frac{1}{3}}(x)} \right] \left[ \frac{\delta \ln q^{\frac{1}{3}}(x)}{\delta h} \right] d^3x = \left[ \int \frac{\bar{H}(x)}{\beta} d^3x \right] \Psi, \quad (13)$$

which describes evolution with respect to intrinsic superspace time interval  $\delta h$ . The corresponding physical Hamiltonian  $\mathcal{H}_{phys.} := \int \frac{\bar{H}(x)}{\beta} d^3x$  is moreover spatial diffeomorphism invariant as it is the integral of a tensor density of weight one. This remarkable Schrodinger equation dictates quantum geometrodynamics in *explicit* superspace  $^{(3)}\mathcal{G}$  entities  $(\Psi[[q_{ij}] \in ^{(3)}\mathcal{G}], \mathcal{H}_{phys.}, \delta h)$ .

On equating  $\delta t = \mathcal{K} \delta h$  in (11), the emergent classical space-time from constructive interference under superspace intrinsic time evolution will, as demonstrated in earlier subsections, be described by the ADM metric,

$$ds^2 = - \left[ \frac{\delta h - \mathcal{L}_{\bar{N}^i} \ln q^{\frac{1}{3}}}{4\beta\kappa(\bar{H}/\sqrt{q})} \right]^2 + q_{ij}[dx^i + \mathcal{N}^i \delta h][dx^j + \mathcal{N}^j \delta h], \quad (14)$$



wherein  $\mathcal{N}^i := N^i/\mathcal{K}$ . Since it can always be absorbed into the gauge parameter  $N^i$ , the constant  $\mathcal{K}$  has no physical implication. Rather, it is the (scalar function)  $\kappa\bar{H}/\sqrt{q}$  which provides the physical conversion between dimensionless intrinsic time interval  $\delta h$  and the proper time of  $ds^2$ . Gravitational redshifts and other physical effects are thus determined by the Hamiltonian density  $\bar{H}/\beta$ ; and the proper time (with vanishing shifts and  $dx^i = 0$ )  $d\tau = \frac{\delta h \sqrt{q}}{4\beta\kappa\bar{H}}$  exhibits the physically intuitive property of varying directly with intrinsic superspace time interval and reciprocally with energy density.

The crucial time development operator can be derived by integrating the Schrodinger equation. This is now feasible without ambiguity because  $\delta h$  is “1-dimensional”, more precisely,  $x$ -independent, rather than many-fingered. Moreover, the necessity of “time”-ordering, which underpins the notion of causality, emerges because quantum fields do not commute at different “times”. Equation (13) implies  $\delta\Psi = [-\frac{i}{\hbar}\mathcal{H}_{phys.}]\delta h\Psi$ ; thus yielding  $\Psi[[q_{ij}(h)] \in {}^{(3)}\mathcal{G}] = U(h, h_0)\Psi[[q_{ij}(h_0)] \in {}^{(3)}\mathcal{G}]$ , with  $h$ -ordered evolution operator  $U(h, h_0) := T \exp \left[ -\frac{i}{\hbar} \int_{h_0}^h \mathcal{H}_{phys.}(h') \delta h' \right]$ . As  $\mathcal{H}_{phys.}$  is classically real and gauge-invariant, a unitary and diffeomorphism-invariant  $U(h, h_0)$  is viable. Since  $\delta h$  is also unchanged under spatial diffeomorphisms, the temporal ordering in  $U(h, h_0)$  is assuredly gauge invariant.

## FURTHER DISCUSSIONS

That there are, for pure gravity, two physical degrees of freedom can be ascertained in the following way: spatial diffeomorphism invariance constraints the physical momenta  $\tilde{\pi}^{ij}$  to be transverse ( $\nabla_i \tilde{\pi}^{ij} = 0$ ) leaving 3 remaining degrees. The 2 transverse traceless modes can be obtained through  $\pi_{TT}^{ij} = (\tilde{\pi}^{ij} - \frac{q^{ij}}{3}\pi) - (\nabla^i W^j + \nabla^j W^i - \frac{2q^{ij}}{3}\nabla_k W^k)$ , with  $W^i$  the solution for  $\nabla_i \pi_{TT}^{ij} = 0$ . Substituting this decomposition of  $\tilde{\pi}^{ij}$  into the symplectic potential, and integrating by parts terms with  $W^i$ , reveal that

$$\int \tilde{\pi}^{ij} \delta q_{ij} = \int \left( \pi_{TT}^{ij} q^{\frac{1}{3}} \delta \bar{q}_{ij} - 2W^j q^{\frac{1}{3}} \nabla^i \delta \bar{q}_{ij} + \pi \delta \ln q^{\frac{1}{3}} \right), \quad (15)$$

which yields  $\int (\tilde{\pi}_T^{ij} \delta \bar{q}_{ij}^{\text{phys.}} + \pi \delta \ln q^{\frac{1}{3}})$  when restricted to the physical subspace with  $\nabla^i \delta \bar{q}_{ij}^{\text{phys.}} = 0$  (this condition has the geometrical meaning physical,  $\delta \bar{q}_{ij}^{\text{phys.}}$ , and gauge,  $\delta \bar{q}_{ij}^{\text{gauge}} = \mathcal{L}_{\bar{N}} \bar{q}_{ij}$ , directions are orthogonal w.r.t. the spatial supermetric  $\bar{G}^{ijkl}$ ). This decomposition yields 2 physical canonical degrees of freedom ( $\bar{q}_{ij}^{\text{phys.}}, \bar{\pi}_T^{ij} := \pi_{TT}^{ij} q^{\frac{1}{3}}$ ), and an extra pair  $(\ln q^{\frac{1}{3}}, \pi)$  to play the role of time and Hamiltonian (which is consistently tied to  $\pi$  by the dynamical equations (3) and (4)). For perturbations about any background  $q_{ij}^* = q_{ij} - \delta q_{ij}$ , the linearized physical spatial metric modes  $\delta \bar{q}_{ij}^{\text{phys.}} = (P_{ij}^{kl})^* \delta q_{kl}$  are traceless ( $q^{*ij} \delta \bar{q}_{ij}^{\text{phys.}} = 0$ ) and transverse ( $\nabla^{*i} \delta \bar{q}_{ij}^{\text{phys.}} = 0$ ) w.r.t  $q_{ij}^*$ , correctly accounting for the perturbative graviton degrees of freedom.

With regard to Lorentz invariance, it should be pointed out although there is only spatial diffeomorphism invariance, for any classical ADM space-time Lorentz symmetry of the tangent space is intact, as the ADM metric  $ds^2 = \eta_{AB} e_\mu^A e_\nu^B dx^\mu dx^\nu = -N^2 dt^2 + q^{\frac{1}{3}} \bar{q}_{ij}(x, t)(dx^i + N^i dt)(dx^j + N^j dt)$  is invariant under local Lorentz transformations of the vierbein fields  $e_\mu^A = \Lambda^A_B(x) e_\mu^B$  which do not affect metric components  $g_{\mu\nu} = \eta_{AB} e_\mu^A e_\nu^B$ . At the more fundamental level of spatial metrics, expressing  $q_{ij} = e_{Ai} e_j^A$  leads to the decomposition of the symplectic potential as

$$\int \tilde{\pi}^{ij} \delta q_{ij} = \int \pi_A^i \delta e_i^A = \int \left( \pi \delta \ln q^{\frac{1}{3}} + \bar{\pi}_A^i \delta \bar{e}_i^A \right); \quad (16)$$

wherein  $\pi^{Ai} := 2\tilde{\pi}^{ij} e_j^A$ ,  $\pi = q_{ij} \tilde{\pi}^{ij} = \frac{1}{2}(\pi_A^i e_i^A)$ ,  $\bar{\pi}^{Ai} := \pi^{Ai} - q^{ij} e_j^A \frac{(\pi_B^k e_k^B)}{3}$ , and  $\bar{e}_i^A := q^{-\frac{1}{6}} e_i^A$ ; together with the Gauss Law constraint,  $\frac{1}{2}[e_i^A \pi^{Bi} - e_i^B \pi^{Ai}] \approx 0$ , which generates  $SO(1, 3)$  Lorentz transformations of the variables.  $(\ln q^{\frac{1}{3}}, \pi)$  is Lorentz invariant and remains decoupled from the rest of the conjugate pairs; and the 2 d.o.f. for pure gravity is maintained.

The inclusion of matter and other forces is rather straightforward as Standard Model fields do not couple to  $\pi$ ; and the corresponding Hamiltonian density of these fields,  $H_M$ , can be appended to gravitational kinetic and potential energy terms. Consequently  $\bar{H}_T = \sqrt{\bar{H}^2 + H_M^2}$  replaces  $\bar{H}$  of (12), and the extension does not affect the fundamental form of the dynamical equation which is now  $\beta\pi + \bar{H}_T = 0$ . In  $H_T$ , Weyl fermions (or primed and unprimed spinors) couple to  $SL(2, C)$  spin connections. For any classical ADM space-time, these spin connections which couple to Standard Model fermions are pullbacks of self- or anti-self-dual Lorentz spin connections to 3-dimensional spatial slices[13]. Isomorphism of  $e_i^A$  to spinorial entities  $e_i^{\mathcal{B}\mathcal{B}'} := e_{Ai}[\tau^A]^{\mathcal{B}\mathcal{B}'}$  is furnished by the quaternion  $\tau^{A=0,1,2,3} := (I_2, \frac{\sigma^{a=1,2,3}}{i})$ .

Identification of a complete set of observables in theories with diffeomorphism invariance is often thought to be more than a challenging task. In this context, for a theory with HJ equation first-order-in-time, the solution is complete[9] in that it has as many integration constants (denoted earlier by  $\alpha$ ) as the number of degrees of freedom in the theory, plus an overall additive constant. These are all gauge invariant, and together with  $\omega := \frac{\delta S}{\delta \alpha}$  (which express the coordinates in terms of time and the constants  $(\alpha, \omega)$ ), provide general integrals of equations of motion which are well-suited to the role physical observables of diffeomorphism-invariant theories. Alternatively, emergent space-time manifolds obtained by integrating Hamilton's equations are characterized by  $4 \times \infty^3$  freely specifiable initial data  $(\bar{q}_{ij}^{\text{phys}}, \bar{\pi}_{TT}^{ij})$ . In a theory with only spatial diffeomorphism gauge symmetry, physical observables are required to commute only with  $H_i$  (and not  $H$ ), thus all Kuchar observables[14] *become* physical. These observables will also consistently have weakly vanishing Poisson bracket with  $M$  (since  $\{f(q_{ij}, \bar{\pi}^{ij}), M\}|_{M=0 \Leftrightarrow H=0} \approx 0$ ).

The symplectic 1-form  $\int \pi \delta \ln q^{\frac{1}{3}} = \int \frac{2}{3}(\frac{\pi}{\sqrt{q}}) \delta \sqrt{q}$  allows a different perspective. With the further restriction of  $\nabla_i \pi = 0 \Leftrightarrow \frac{\pi}{\sqrt{q}} = T$ , York[12] interprets and deploys the scalar  $\frac{\pi}{\sqrt{q}} = -\frac{TK}{6\beta^2\kappa}$  as the “extrinsic time” variable. It follows that the Hamiltonian  $\bar{H} = -\beta\pi$  is then proportional to  $\sqrt{q}$ , and the total energy to the volume. In our framework, York's restriction of spatially constant  $\frac{\pi}{\sqrt{q}} = -\frac{\bar{H}}{\beta\sqrt{q}} = T$  is a special case wherein, with vanishing shift vectors,  $d\tau^2 = [\frac{\delta h}{4\beta^2\kappa T}]^2$ . Although the extrinsic time variable is then invariant under spatial diffeomorphisms, it is however not invariant under 4-dimensional coordinate transformations which are supposedly symmetries of Einstein's theory. With the paradigm shift to just spatial diffeomorphism invariance,  $\delta \ln q^{\frac{1}{3}}$  is well-suited to the role of physical time interval: it is a spatial diffeomorphism scalar with a gauge-invariant part  $\delta h$  which is spatially constant.

From the Schrodinger and HJ equations, pure General Relativity has the physical content of conjugate variables  $(\bar{q}_{ij}, \bar{\pi}^{ij})$  subject to  $H_i = 0$  evolving w.r.t.  $\ln q^{\frac{1}{3}}$  with effective Hamiltonian density  $\bar{H}/\beta$ . Thus proceeding from the action  $\int [\bar{\pi}^{ij} \delta \bar{q}_{ij} - \frac{\bar{H}}{\beta} \delta \ln q^{\frac{1}{3}}] - \int \delta \mathcal{N}^i H_i$ , and inverting for  $\bar{\pi}^{ij}$  in terms of  $\frac{\delta \bar{q}_{ij}}{\delta \ln q}$  from the EOM, yields (the result can also be deduced from (5) and (6)) the action functional as

$$S = - \int \sqrt{V} \sqrt{\frac{1}{\beta^2} (\delta \ln q^{\frac{1}{3}} - \mathcal{L}_{\delta \bar{N}} \ln q^{\frac{1}{3}})^2 - \bar{G}^{ijkl} (\delta \bar{q}_{ij} - \mathcal{L}_{\delta \bar{N}} \bar{q}_{ij}) (\delta \bar{q}_{kl} - \mathcal{L}_{\delta \bar{N}} \bar{q}_{kl})}, \quad (17)$$

which is just the superspace proper time with  $\sqrt{V}$  playing the role of “mass” if it were constant. This regains the generalized Baierlein-Sharp-Wheeler action[15] which has also been studied in Ref.[16] in a different situation.

Transparent and consistent dynamics revealing an emergent lapse, the primacy of the physical Hamiltonian  $\mathcal{H}_{\text{phys.}}$ , and the role of intrinsic time in General Relativity[17] and its extensions can be obtained from several complementary approaches: the master constraint formulation which recovers the correct physical content, without paradigm of space-time covariance, from the usual starting point of canonical General Relativity; the Schrodinger equation (3) (or its superspace version (13)) as the fundamental equation for quantum geometrodynamics; the generalized Baierlein-Sharp-Wheeler action (17); and also, perhaps most important to a causal quantum theory, the evolution operator  $U(h, h_0)$  with gauge-invariant temporal ordering.

## APPENDIX

The simple relativistic point particle analogy is quite instructive. The action is

$$\begin{aligned} S &= -m_0 c \int \sqrt{-\eta_{\mu\nu} dx^\mu dx^\nu} = -m_0 c^2 \int d\tau \\ &= -m_0 c^2 \int dt \sqrt{1 - (\frac{d\vec{x}}{cdt})^2} = \int (\vec{p} \cdot \frac{d\vec{x}}{dt} - c \sqrt{\vec{p}^2 + m_0^2 c^2}) dt, \end{aligned} \quad (18)$$

with  $\vec{p} = m_0 \frac{d\vec{x}}{d\tau}$ , and the physical Hamiltonian  $\bar{H} = c \sqrt{\vec{p}^2 + m_0^2 c^2}$  emerges when  $t$  is correctly identified as the time variable. On the other hand, in the “manifestly covariant” approach, introduction of an extraneous “time” parameter  $\lambda$  results in

$$S = -m_0 c \int \sqrt{-\eta_{\mu\nu} \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}} d\lambda = \int \mathcal{L}[x^\mu, \frac{dx^\mu}{d\lambda}] d\lambda \quad \Rightarrow \quad p_\mu = \frac{\partial \mathcal{L}}{\partial (\frac{dx^\mu}{d\lambda})} = m_0 \eta_{\mu\nu} \frac{dx^\nu}{d\tau},$$

with  $p_\mu \frac{dx^\mu}{d\lambda} - \mathcal{L} = 0$  i.e. exactly vanishing Hamiltonian; but the momenta are constrained by  $H = p^\mu p_\mu + m_0^2 c^2 = -(p_0 - \frac{\bar{H}}{c})(p_0 + \frac{\bar{H}}{c}) = 0$ . Formulating the constrained theory with

$$S = \int [p_\mu \frac{dx^\mu}{d\lambda} - N(p^\mu p_\mu + m_0^2 c^2)] d\lambda \quad (19)$$



results in EOM which (a posteriori) determine  $Nd\lambda = \frac{cdt}{2p^0} = \frac{cdt}{(2\sqrt{\vec{p}^2 + m_0^2 c^2})} = \frac{c^2 dt}{2\bar{H}}$ . The obfuscating reparametrization “symmetry” associated with  $\lambda$  (which is not intrinsic to the theory) and constrained Hamiltonian  $N(p^\mu p_\mu + m_0^2 c^2)$  gives rise to artificial gauge histories of  $x^\mu$  in  $\lambda$ -time with Lagrange multiplier  $N$ . Correctly isolating one of the degree  $x^0$  as the intrinsic time, and forgoing the “ $\lambda$ -reparametrization symmetry” regains the much more transparent description of (18) with dynamical variables  $\vec{x}$  evolving w.r.t. intrinsic time  $t = \frac{x^0}{c}$  and physical Hamiltonian  $\bar{H}$ . In the event (20) is the starting point (analogous to the situation in General Relativity), the physical description of (18) can be recovered by introducing the master constraint  $M = (p_0 + \frac{\bar{H}}{c})^2 = 0$  (which makes  $H = 0$  redundant and in effect replaces it); thus yielding

$$\begin{aligned} S &= \int [p_\mu \frac{dx^\mu}{d\lambda}] d\lambda - \int m(p_0 + \frac{\bar{H}}{c})^2 d\lambda = \int [p_0 c + \vec{p} \cdot \frac{d\vec{x}}{dt}] dt - \int m(p_0 + \sqrt{\vec{p}^2 + m_0^2 c^2})^2 d\lambda \\ &= \int [-c\sqrt{\vec{p}^2 + m_0^2 c^2} + \vec{p} \cdot \frac{d\vec{x}}{dt}] dt + \int m(p_0 + \sqrt{\vec{p}^2 + m_0^2 c^2})^2 d\lambda, \end{aligned} \quad (20)$$

which implies  $\bar{H} = c\sqrt{\vec{p}^2 + m_0^2 c^2}$  is the effective Hamiltonian for the variables  $(\vec{x}, \vec{p})$ , and the  $M$  constraint is equivalent to  $p_0 = -\frac{\bar{H}}{c}$  which can consistently be interpreted classically as the Hamilton-Jacobi equation  $\frac{\partial S}{\partial x^0} + \frac{\bar{H}}{c} = 0$ , and quantum mechanically as a Schrodinger equation.

## ACKNOWLEDGMENTS

This work has been supported in part by the National Science Council of Taiwan under Grant Nos. NSC98-2112-M-006-006-MY3, 101-2112-M-006-007-MY3; the Institute of Physics, Academia Sinica; and the National Center for Theoretical Sciences, Taiwan.

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